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## Mn<sup>II</sup>[Mn<sup>II</sup>(CN)<sub>4</sub>]<sub>2</sub>—A Magnetic Interpenetrating Three-Dimensional Diamondlike Solid\*\*

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The preparation of the first paramagnetic tetrahedral tetracyanometalate<sup>[1]</sup> enables the construction new magnetic 3-D network solids, as [Mn<sup>II</sup>(CN)<sub>4</sub>]<sup>2-</sup> is a building block for the preparation of four-coordinate extended network structures akin to the interpenetrating diamondlike Zn(CN)<sub>2</sub>.<sup>[2]</sup> Information on the magnetic properties of this new class of materials will enable a direct comparison to the six-coordinate Prussian Blue family of ordered magnets.<sup>[3]</sup>

The reaction of equimolar amounts of [PPN]<sub>2</sub>[Mn<sup>II</sup>(CN)<sub>4</sub>] (PPN = [Ph<sub>3</sub>P=N=PPh<sub>3</sub>]<sup>+</sup>) and [Mn<sup>II</sup>(NCMe)<sub>6</sub>][B(3,5-C<sub>6</sub>H<sub>3</sub>(CF<sub>3</sub>)<sub>2</sub>)<sub>4</sub>] leads to the immediate precipitation of a solid of composition Mn(CN)<sub>2</sub>, formulated as Mn<sup>II</sup>·[Mn<sup>II</sup>(CN)<sub>4</sub>]. This red solid exhibits an  $\tilde{\nu}_{\text{CN}}$  absorption at 2170 cm<sup>-1</sup>, shifted relative to the corresponding band of [Mn<sup>II</sup>(CN)<sub>4</sub>]<sup>2-</sup> by 32 cm<sup>-1</sup> to lower energy, which suggests the lengthening of the CN bond upon coordination of its N atom to high-spin Mn<sup>II</sup>. Based on the 2 $\theta$  values and relative intensities, the powder pattern is assigned to the cubic *P*43*m* space group with  $a = 6.1229(5)$  Å (Figure 1).<sup>[4]</sup> This is consistent with an interpenetrating diamondlike (sphalerite) lattice (Figure 2), based on tetrahedral Mn<sup>II</sup> sites with an Mn<sup>II</sup>...Mn<sup>II</sup> separation of 5.303 Å, which is 0.20 Å longer than the intermetallic distance found for

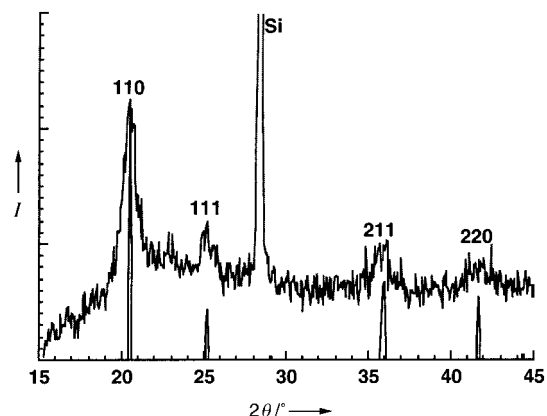


Figure 1. Observed (top) and calculated (lower) powder X-ray diffraction pattern for Mn<sup>II</sup>[Mn<sup>II</sup>(CN)<sub>4</sub>]<sub>2</sub>. The calculated pattern is for the *P*43*m* space group with  $a = 6.1229(5)$  Å. Silicon was used as an internal standard.  $I$  = intensity.

Zn(CN)<sub>2</sub> and 0.23 Å longer than that in K<sub>2</sub>Mn<sup>II</sup>[Mn<sup>II</sup>(CN)<sub>6</sub>].<sup>[7]</sup> Solvent is not evident from either IR or thermal gravimetric analyses. The interpenetrating structure would lead to a density twice the value of 0.775 g cm<sup>-3</sup> (i.e., 1.550 g cm<sup>-3</sup>) calculated for a noninterpenetrating structure; indeed, 1.68 g cm<sup>-3</sup> is observed for a small sample.<sup>[8]</sup> Due to the inherent lability and coordinative unsaturation of the 13 e<sup>-</sup>

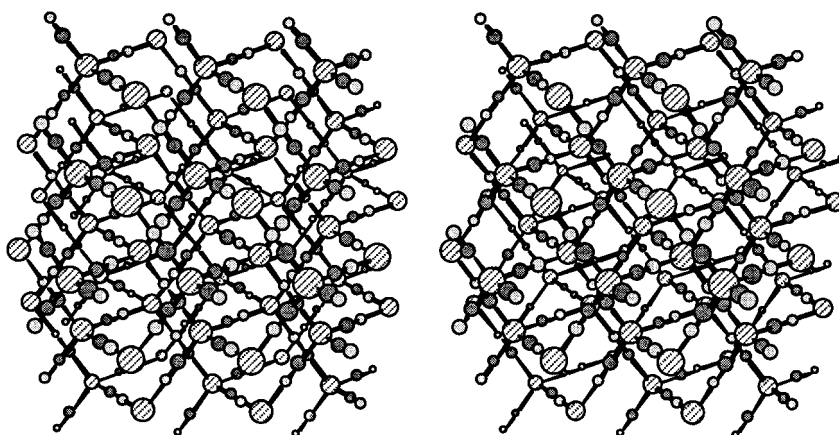


Figure 2. Stereoview of the proposed double-density interpenetrating diamondlike lattice for Mn<sup>II</sup>[Mn<sup>II</sup>(CN)<sub>4</sub>]<sub>2</sub>. The structure resembles that of sphalerite.

[Mn(CN)<sub>4</sub>]<sup>2-</sup> building block in polar organic solvents, the CN ligand can dissociate or bind to other [Mn(CN)<sub>4</sub>]<sup>2-</sup> ions permitting the competitive formation of by-products such as octahedral complexes. Consequently, some variability in the energy of the  $\tilde{\nu}_{\text{CN}}$  band and in the magnetic properties is observed.

The room-temperature magnetic moment  $\mu$  is 4.4  $\mu_{\text{B}}$ , a value substantially lower than the predicted value (5.92  $\mu_{\text{B}}$ ) for uncoupled Mn<sup>II</sup> ions with  $S = 5/2$ , but consistent with strong antiferromagnetic coupling. As depicted in Figure 3 the magnetic susceptibility  $\chi$  gradually increases with decreasing temperature reaching a broad maximum, and then decreases rapidly upon further cooling to 2 K. The broad maximum is consistent with antiferromagnetic ordering, while the upturn at low temperature is assigned to impurities with  $S = 5/2$ . Mn<sup>II</sup>

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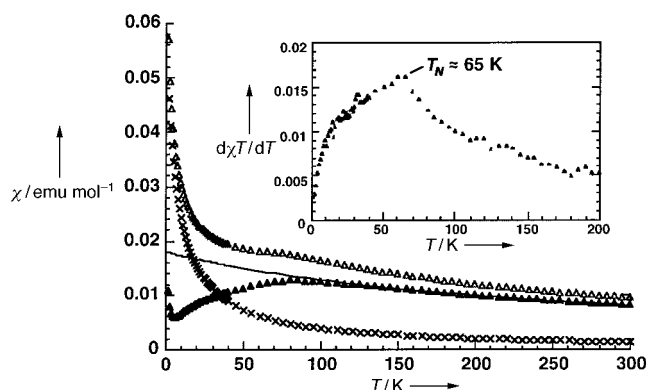


Figure 3. Uncorrected ( $\Delta$ ) and corrected ( $\blacktriangle$ ) molar magnetic susceptibility  $\chi$  as a function of temperature  $T$  for  $\text{Mn}^{\text{II}}[\text{Mn}^{\text{II}}(\text{CN})_4]$ , as well as that obtained after correction for a spin impurity (9.5%) containing  $\text{Mn}^{\text{II}}$  with  $S = 5/2$  ( $\times$ ). The fit above 140 K to the Curie–Weiss expression is also presented (—). The inset shows the maximum at  $T_N \approx 65$  K from the  $d\chi(T)/dT$  ( $\blacktriangle$ ) data.

spin impurities are believed to arise as a result of simultaneous formation of an insoluble oligomeric species as the solid precipitates from solution. Subtraction of a spin impurity<sup>[6]</sup> leads to the corrected susceptibility with a broad maximum at 80 K due to 3-D antiferromagnetic ordering (Figure 3). Above 140 K  $\chi$  can be fitted by the Curie–Weiss expression,  $\chi \propto (T - \theta)^{-1}$  with a strongly antiferromagnetic  $\theta$  value of  $-240$  K (Figure 4). The Neél ordering temperature  $T_N$  is  $\sim 65$  K, as determined from the maximum in the  $d\chi/dT$  data (Figure 3 inset). Alternating current susceptibility  $\chi_{ac}$  below 75 K further substantiates antiferromagnetic ordering, as a peak in  $\chi''$  is not observed.

Although both  $\text{K}_2\text{Mn}^{\text{II}}[\text{Mn}^{\text{II}}(\text{CN})_6]$ <sup>[3c]</sup> and  $\text{Mn}^{\text{II}}[\text{Mn}^{\text{II}}(\text{CN})_4]$  have  $\text{Mn}^{\text{II}} \leftarrow \text{C} \equiv \text{N} \rightarrow \text{Mn}^{\text{II}}$  linkages (six for the former and four for the latter), the former has an uncompensated moment and is a ferrimagnet below 41 K, whereas the latter has a compensated moment and is an antiferromagnet with a higher ordering temperature. This is a consequence of the different electronic structures. Each substance has high-spin  $\text{Mn}^{\text{II}}$  ( $S = 5/2$ ) centers bound to the N atoms. However, the six C-bound CN ligands provide a strong field and the  $\text{Mn}^{\text{II}}$  bound to them is low-spin ( $S = 1/2$ ) and in contrast, as noted above, the tetrahedral  $\text{Mn}^{\text{II}}$  bound to C atoms of the CN ligands is high-spin ( $S = 5/2$ ). Hence,  $\text{K}_2\text{Mn}^{\text{II}}[\text{Mn}^{\text{II}}(\text{CN})_6]$  has

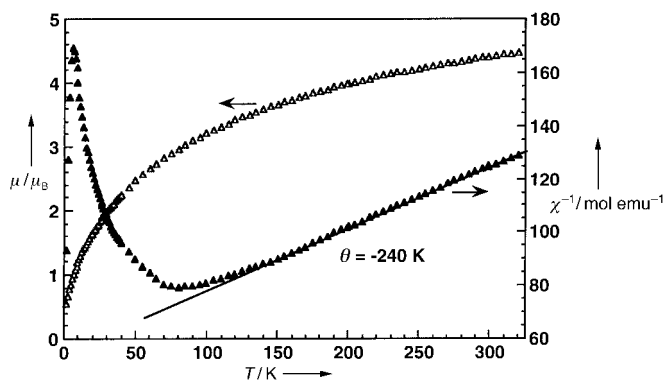


Figure 4. Corrected reciprocal molar magnetic susceptibility  $\chi^{-1}$  ( $\blacktriangle$ ) and moment  $\mu/\mu_B$  ( $\Delta$ ) as a function of temperature  $T$  for  $\text{Mn}^{\text{II}}[\text{Mn}^{\text{II}}(\text{CN})_4]$ . Above 140 K  $\chi^{-1}(T)$  is linear and can be fitted by the Curie–Weiss expression with  $\theta = -240$  K.

coupling between the high- and low-spin sites ( $S = 5/2 - 1/2 = 2$ ) leading to an uncompensated moment, while both  $\text{Mn}^{\text{II}}$  centers have a spin  $S = 5/2$  for  $\text{Mn}^{\text{II}}[\text{Mn}^{\text{II}}(\text{CN})_4]$  and lead to a compensated moment, that is,  $S = 5/2 - 5/2 = 0$ .

$T_c$  can be estimated<sup>[3b, f]</sup> from Equation (a), where  $J$  is the nearest neighbor magnetic coupling constant, and  $C_{\text{MC}}$  and  $C_{\text{MN}}$  are the Curie constants for the C- and N-bound  $\text{Mn}^{\text{II}}$

$$T_c \propto zJ[C_{\text{MC}}C_{\text{MN}}]^{1/2} \quad (\text{a})$$

centers. For systems with a Landé factor  $g$  of 2,  $C$  is proportional to  $S(S+1)$ , and assuming identical  $J$  values the  $T_c$  for tetrahedral  $\text{Mn}^{\text{II}}[\text{Mn}^{\text{II}}(\text{CN})_4]$  should be 2.28 times greater than the 41 K (i.e., 93.5 K), which is observed for octahedral  $\text{K}_2\text{Mn}^{\text{II}}[\text{Mn}^{\text{II}}(\text{CN})_6]$ . Since only 1.58 times the 41 K is observed,  $J$  for  $\text{Mn}^{\text{II}}[\text{Mn}^{\text{II}}(\text{CN})_4]$  must be smaller. This is consistent with the substantially longer  $\text{Mn}^{\text{II}}-\text{C}$  bond (by about 0.20 Å), which reduces  $J$ . An independent estimation of  $J$  requires a model for a 3-D diamond-type Heisenberg system with  $S = 5/2$ , which at present does not exist.

In contrast to  $\text{Mn}^{\text{II}}[\text{Mn}^{\text{II}}(\text{CN})_4]$ ,  $\text{M}'[\text{Mn}^{\text{II}}(\text{CN})_4]$  ( $\text{M} \neq \text{Mn}$ ) should have uncompensated moments and order ferrimagnetically. Studies on these materials are in progress.

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- [4] Based upon the corrected (with respect to Si internal standard)  $d$ -spacings, a crude lattice parameter was calculated and peak indices using  $\text{Zn}(\text{CN})_2$  as a structural model.<sup>[2]</sup> Using LATPARM and assuming cubic symmetry the reflection indices along with  $d$ -spacings were refined. The refinement converged after only four cycles and generated a new lattice parameter with the standard deviation. The space group  $P43m$  assumes an ordered arrangement of cyano ligands that is  $\text{MnC}_4$  and  $\text{MnN}_4$ , which in reality may not be reasonable. Based on a random orientation of the cyano ligands, that is,  $\text{ZnCN}_3$ ,  $\text{ZnC}_2\text{N}_2$ ,  $\text{ZnC}_3\text{N}$ ,  $\text{ZnC}_4$ , and  $\text{ZnN}_4$ , a lower symmetry space group,  $Pn3m$ , has been proposed.<sup>[5]</sup> Altering the space group merely changes (slightly) the calculated peak intensities. A calculated X-ray diffraction pattern was generated with CRYSTALDIFFRACT<sup>[6]</sup> on using the refined lattice parameter and atomic positions of the isostructural  $\text{Zn}(\text{CN})_2$  compound. Therefore the calculated and experimental diffraction patterns are in excellent agreement.
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- [7] The  $\text{Mn}^{\text{II}} \cdots \text{Mn}^{\text{II}}$  separation is determined from the positions of the Mn centers (Mn Wyckoff positions: 1a  $\text{Mn}(1)$  0,0,0; 3d  $\text{Mn}(2)$  0,0, $1/2$  as follows:  $3^{1/2}a/2$ .
- [8] The density of the material was determined by utilizing a  $\text{CH}_2\text{I}_2/\text{CH}_2\text{Cl}_2$  mixture and weighing a 1.00 mL volume.
- [9] The impurity (9.5%) is not a simple paramagnet, but possesses some antiferromagnetic coupling ( $\theta = -9$  K). This coupling suggests nearest neighbor interactions among adjacent  $\text{Mn}^{\text{II}}$  sites.